

The 4th workshop on spectral graph theory and related topics  
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## Complex Fermi surfaces and spectrum of discrete Laplacian on perturbed lattices

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# Contents

I'll talk about eigenvalues embedded in the continuous spectrum of discrete Laplacian on lattices with finite rank perturbations :

- ▶ discrete Laplacian on periodic lattices
- ▶ Fermi surfaces
- ▶ unique continuation properties and lattice structures
- ▶ Rellich type uniqueness theorem and its applications

## Recent progress on discrete Laplacian

- ▶ Spectrum of Laplacian on infinite graphs (Higuchi-Shirai, 2004 / Higuchi-Nomura, 2009)
- ▶ Endpoint embedded eigenvalue for higher dim. (Hiroshima-Sakai-Sasaki-Suzuki, 2012)
- ▶ Absence of embedded eigenvalues on the square lattice (Isozaki-Morioka, 2014)
- ▶ Generalization for some kind of periodic lattices (Ando-Isozaki-Morioka, 2015)
- ▶ Endpoint embedded eigenvalue for low dim. (Ogurusu-Higuchi-Nomura, in preparation?)
- ▶ Generalization for exponential decaying perturbations (Vesalainen, 2014)
- ▶ Generalizations for short-range perturbations (Morioka, in preparation)
- ▶ Tree (Colin de Verdière-Truc, 2013)
- ▶ Periodic lattices with pendant vertices (Suzuki, 2013)

## Discrete Laplacian with finite rank perturbations

- ▶ H. Isozaki, H. Morioka, Inverse Problems and Imaging, **8** (2014), 475-489.
- ▶ K. Ando, H. Isozaki, H. Morioka, Ann. Henri Poincaré, online first (2015).

## Periodic lattices

## e.g.: Square lattice

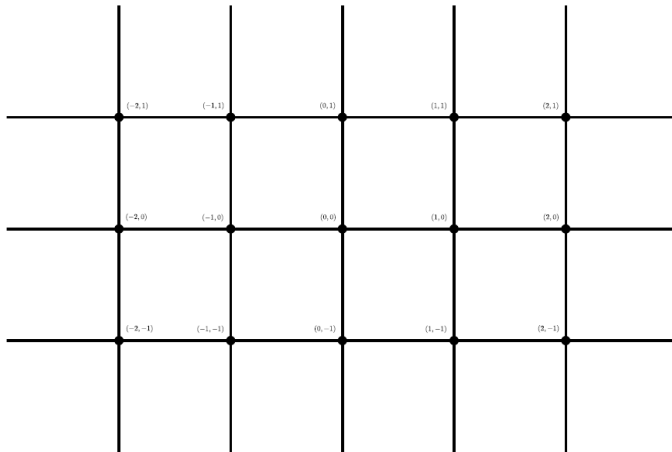


Figure: Square lattice

## e.g.: Hexagonal and diamond lattice

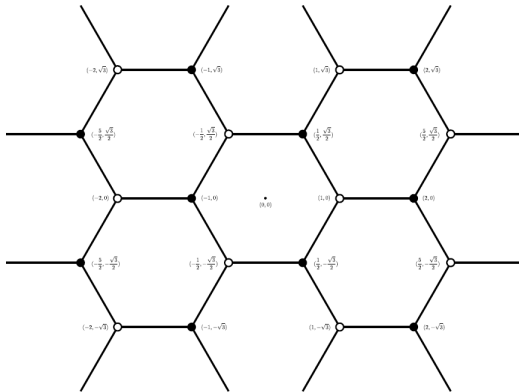


Figure: Hexagonal lattice

## e.g.: Kagome lattice

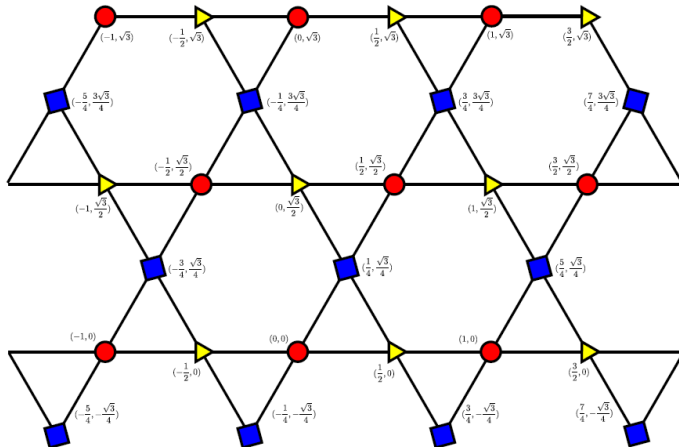


Figure: Kagome lattice



e.g.: Ladder

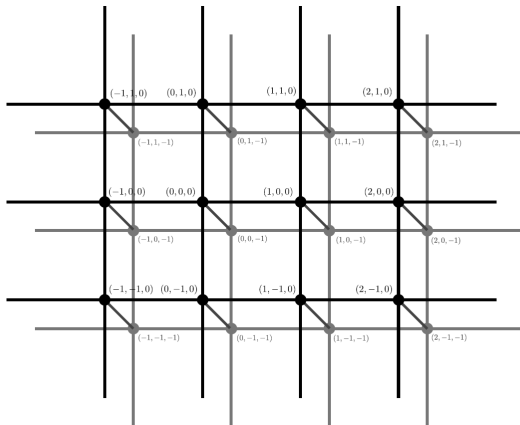


Figure: Ladder of  $\mathbb{Z}^2$

e.g.: Graphite

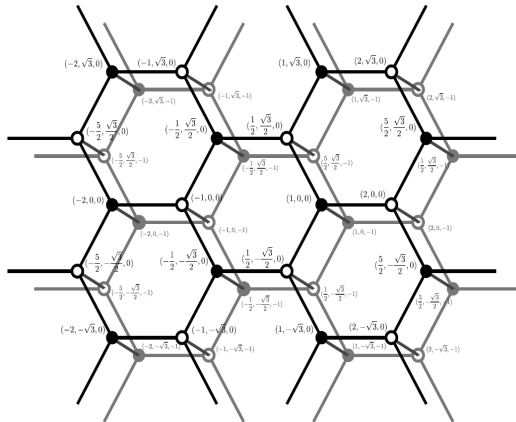


Figure: Graphite

## e.g.: subdivision of the 2-D square lattice

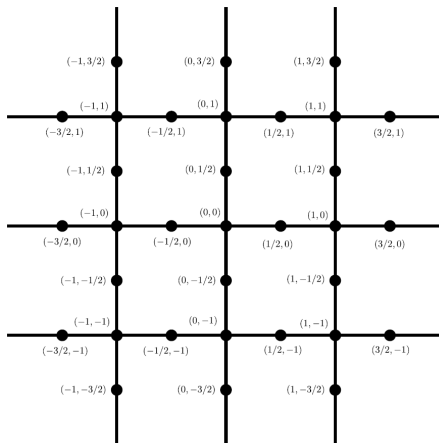


Figure: subdivision

## Definition of lattices

These lattice are constructed as a  $\mathbf{Z}^d$ -covering of a finite graph :

We define  $\Gamma_0 = \{\mathcal{L}_0, \mathcal{V}_0, \mathcal{E}_0\}$  by :

For the basis  $\mathbf{v}_j, j = 1, \dots, d$  of  $\mathbf{R}^d$ , we put

$$\mathcal{L}_0 = \{\mathbf{v}(n) ; n \in \mathbf{Z}^d\}, \quad \mathbf{v}(n) = \sum_{j=1}^d n_j \mathbf{v}_j, \quad n \in \mathbf{Z}^d.$$

and, for some points  $p_1, \dots, p_s$  in  $\mathbf{R}^d$ , we define the set of vertices by

$$\mathcal{V}_0 = \bigcup_{j=1}^s (p_j + \mathcal{L}).$$

Moreover, we assume that the set of unoriented edges  $\mathcal{E}_0$  is invariant with respect to  $\mathbf{Z}^d$ -action.

We assume that  $\Gamma_0$  has no self-loops nor multiple edges.

## Discrete Laplacian (transition Laplacian)

For a  $\mathbf{C}$ -valued function  $\widehat{f} = \{\widehat{f}(v)\}_{v \in \mathcal{V}_0}$ , we define

$$(\widehat{\Delta}_{\Gamma_0} \widehat{f})(v) = \frac{1}{\deg(v)} \sum_{w \in \mathcal{V}_0, (v,w) \in \mathcal{E}_0} \widehat{f}(w),$$

$$\deg(v) := \sharp\{w \in \mathcal{V}_0 ; (v, w) \in \mathcal{E}_0\}.$$

By a  $\mathbf{R}$ -valued scalar potential  $\widehat{V}$ , we define the discrete Schrödinger equation by

$$(-\widehat{\Delta}_{\Gamma_0} + \widehat{V} - \lambda)\widehat{u} = 0 \quad \text{on} \quad \mathcal{V}_0.$$

## Systems of Schrödinger equation

In view of periodic structure of  $\Gamma_0$ , we interpret  $\ell^2(\mathcal{V}_0)$  as  $\bigoplus_{j=1}^s \ell^2(\mathbb{Z}^d) = \ell^2(\mathbb{Z}^d; \mathbb{C}^s)$ . (Other function spaces are also dealt with spaces of  $\mathbb{C}^s$ -valued functions on  $\mathbb{Z}^d$ .)

Then we rewrite  $\widehat{f} = \{\widehat{f}(v)\}_{v \in \mathcal{V}_0}$  by

$$\widehat{f}(n) = (\widehat{f}_1(n), \dots, \widehat{f}_s(n)), \quad n \in \mathbb{Z}^d,$$

so that, by the Shift operator  $(\widehat{S}_j^\pm \widehat{f})(n) = \widehat{f}(n \pm e_j)$ , we have

$$\widehat{H}_0 = -\widehat{\Delta}_{\Gamma_0} = s \times s \text{ symmetric matrix of } \widehat{S}_j^\pm,$$

$$\widehat{V} = \text{diag}(\widehat{V}_1, \dots, \widehat{V}_s).$$

$\widehat{H} = \widehat{H}_0 + \widehat{V}$  is bounded and self-adjoint on  $\ell^2(\mathbb{Z}^d; \mathbb{C}^s)$ .

## Fourier transformation

On  $\Gamma_0$ , we can use the Fourier series :

The unitary operator  $\mathcal{U}_{\mathcal{L}_0} : \ell^2(\mathbb{Z}^d; \mathbb{C}^s) \rightarrow L^2(\mathbb{T}^d; \mathbb{C}^s)$  is defined by

$$(\mathcal{U}_{\mathcal{L}_0} \widehat{f})_j(x) = (2\pi)^{-d/2} \sqrt{\deg_0(j)} \sum_{n \in \mathbb{Z}^d} \widehat{f}_j(n) e^{in \cdot x}, \quad j = 1, \dots, s,$$

where  $\deg_0(j)$  is the degree of vertex in each orbit, and the inner products are

$$(\widehat{f}, \widehat{g})_{\ell^2(\mathbb{Z}^d; \mathbb{C}^s)} = \sum_{j=1}^s \sum_{n \in \mathbb{Z}^d} \deg_0(j) \widehat{f}_j(n) \overline{\widehat{g}_j(n)},$$

$$(u, v)_{L^2(\mathbb{T}^d; \mathbb{C}^s)} = \sum_{j=1}^s \int_{\mathbb{T}^d} u_j(x) \overline{v_j(x)} dx.$$

## Fourier transform of Discrete Laplacian

In the following, we denote  $H = H_0 + V$ ,

$$H_0 = \mathcal{U}_{\mathcal{L}_0} \widehat{H}_0 \mathcal{U}_{\mathcal{L}_0}^* = H_0(x),$$

$H_0(x) = s \times s$  Hermitian matrix with trigonometric-function-entries,

$$V = \mathcal{U}_{\mathcal{L}_0} \widehat{V} \mathcal{U}_{\mathcal{L}_0}^*,$$

on the torus  $\mathbf{T}^d$ .

Then the discrete Schrödinger equation on  $\mathbf{T}^d$  is

$$(H_0(x) - \lambda)u + Vu = 0 \quad \text{on} \quad \mathbf{T}^d.$$



## Diagonalization of $H_0(x) - \lambda$

Multiplying the co-factor of  $H_0(x) - \lambda$ , we rewrite  $H_0(x) - \lambda$  as

$$H_0(x) - \lambda \rightarrow p(x, \lambda)I_s, \quad p(x, \lambda) := \det(H_0(x) - \lambda).$$

Putting eigenvalues of  $H_0(x)$  for each  $x \in \mathbb{T}^d$  as  $\lambda_1(x) \leq \cdots \leq \lambda_s(x)$ , we have

$$p(x, \lambda) = \prod_{j=1}^s (\lambda_j(x) - \lambda).$$

## Essential spectrum of $\widehat{H}_0$

## Essential spectrum of $\widehat{H}_0$

The spectrum of  $\widehat{H}_0$  is given by

$$\sigma(\widehat{H}_0) = \sigma(H_0) = \bigcup_{x \in \mathbb{T}^d} \text{“eigenvalue of } H_0(x)\text{”}.$$

- ▶ In our cases,  $\sigma(\widehat{H}_0) = \sigma_{ess}(\widehat{H}_0)$  is a closed interval  $\subset [-1, 1]$ .
- ▶ There may exist  $\lambda \in \sigma_p(\widehat{H}_0) \cap \sigma_{ess}(\widehat{H}_0)$ ! (e.g. Kagome lattice, subdivision lattices)
- ▶ Generally,  $\sigma(\widehat{H}_0)$  may have some spectral gaps, and eigenvalues with  $\infty$ -multiplicities. (c.f. Suzuki, 2013 et al.)

## Examples (1)

(1) : Square lattice

$$\blacktriangleright p(x, \lambda) = -\frac{1}{d} \sum_{j=1}^d \cos x_j - \lambda$$

$$\blacktriangleright \sigma(\widehat{H}_0) = [-1, 1]$$

(2) : Hexagonal lattice

$$\blacktriangleright p(x, \lambda) = \lambda^2 - \frac{3 + 2(\cos x_1 + \cos x_2 + \cos(x_1 - x_2))}{9}$$

$$\blacktriangleright \sigma(\widehat{H}_0) = [-1, 1]$$

(3) : Kagome lattice

$$\blacktriangleright p(x, \lambda) = -\left(\lambda - \frac{1}{2}\right) \left( \lambda^2 + \frac{\lambda}{2} - \frac{1 + \cos x_1 + \cos x_2 + \cos(x_1 - x_2)}{8} \right)$$

$$\blacktriangleright \sigma(\widehat{H}_0) = [-1, 1/2]$$

## Examples (2)

(4) : Subdivision lattices

$$\blacktriangleright p(x, \lambda) = (-\lambda)^{d-1} \left( \lambda^2 - \frac{1}{2d} (d + \sum_{j=1}^d \cos x_j) \right)$$

$$\blacktriangleright \sigma(\widehat{H}_0) = [-1, 1]$$

(5) : (Higher dimensional) ladders

$$\blacktriangleright p(x, \lambda) = p_+(x, \lambda) p_-(x, \lambda)$$

$$\blacktriangleright p_{\pm}(x, \lambda) = \lambda + \frac{1}{2d+1} \left( 2 \sum_{j=1}^d \cos x_j \pm 1 \right)$$

$$\blacktriangleright \sigma(\widehat{H}_0) = [-1, 1]$$

(6) : Graphite

$$\blacktriangleright p(x, \lambda) = \lambda^4 - \frac{\alpha(x) + 1}{8} \lambda^2 + \frac{(\alpha(x) - 1)^2}{4^4}$$

$$\blacktriangleright \alpha(x) = 3 + 2(\cos x_1 + \cos x_2 + \cos(x_1 - x_2))$$

$$\blacktriangleright \sigma(\widehat{H}_0) = [-1, 1]$$

## Remarks

The spectral measure of  $\sigma(\widehat{H}_0)$  highly depends on the structure of lattices :

- ▶ Geometric structure of complex Fermi surfaces of  $\widehat{H}_0$
- ▶ Density of states

## Fermi surfaces

## Real Fermi surfaces, complex Fermi surfaces

The set  $M_\lambda \subset \mathbb{T}^d$  is defined by

$$M_\lambda = \{x \in \mathbb{T}^d ; p(x, \lambda) = 0\}.$$

We extend  $M_\lambda$  to the complex torus  $\mathbb{T}_\mathbb{C}^d = \mathbb{C}^d / (2\pi\mathbb{Z})^d$ , and denote it by

$$M_\lambda^\mathbb{C} = \{z \in \mathbb{T}_\mathbb{C}^d ; p(z, \lambda) = 0\}.$$

We split  $M_\lambda^\mathbb{C}$  into two part, one is the regular part and another is the singular part :

$$M_{\lambda, reg}^\mathbb{C} = \{z \in M_\lambda^\mathbb{C} ; \nabla_z p(z, \lambda) \neq 0\},$$

$$M_{\lambda, sing}^\mathbb{C} = \{z \in M_\lambda^\mathbb{C} ; \nabla_z p(z, \lambda) = 0\}.$$



## Assumption for the Fermi surfaces

### Assumption

(A-1) There exists a subset  $\mathcal{T}_1 \subset \sigma(\widehat{H}_0)$  such that for  $\lambda \in \sigma(\widehat{H}_0) \setminus \mathcal{T}_1$  :

(A-1-1)  $M_{\lambda, \text{sing}}^C$  is a discrete set.

(A-1-2) Each connected component of  $M_{\lambda, \text{reg}}^C$  intersects with  $\mathbf{T}^d$ . Each intersection is a  $(d - 1)$ -dimensional real analytic submanifold of  $\mathbf{T}^d$ .

※ Since  $M_{\lambda}^C$  is defined by the trigonometric polynomial  $p(z, \lambda)$ , we can not define “irreducible factor”. However, we can consider an irreducibility in view of the connectivity as complex submanifolds.

## Rellich type uniqueness theorem

## Rellich type theorem

### Theorem (Ando-Isozaki-Morioka, 2015)

We assume (A-1), and let  $\lambda \in \sigma(\widehat{H}_0) \setminus \mathcal{T}_1$ .

For a function  $f$  whose entries are trigonometric polynomials, suppose a distribution  $u$  satisfies the equation

$$(H_0(x) - \lambda)u = f \quad \text{on } \mathbb{T}^d,$$

$$\lim_{R \rightarrow \infty} \frac{1}{R} \sum_{j=1}^s \int_{\mathbb{T}^d} |\chi(|\sqrt{-\Delta}| < R) u_j(x)|^2 dx = 0.$$

Then entries of  $u$  are also trigonometric polynomials.

※ Recalling  $\widehat{f} := \mathcal{U}_{\mathcal{L}_0} f$ , we have  $\sharp \text{supp } \widehat{f} < \infty$ . Vesalainen (2014) has generalized our result (Isozaki-Morioka, 2014) for infinite rank perturbations with the condition

$$e^{\gamma \langle n \rangle} \widehat{f} \in \ell^2(\mathbb{Z}^d) \text{ for } \forall \gamma > 0, \quad \widehat{f}(n) = 0 \text{ for } \sum_{j=1}^{d-1} |n_j| \leq n_d.$$

## Interpretation on the lattice

### Corollary

We assume (A-1).

If, for a constant  $R_0 > 0$  and  $\lambda \in \sigma(\widehat{H}_0) \setminus \mathcal{T}_1$ ,  $\widehat{u}$  satisfies the equation

$$(-\widehat{\Delta}_{\Gamma_0} - \lambda)\widehat{u} = 0 \quad \text{in} \quad |n| > R_0,$$

$$\lim_{R \rightarrow \infty} \frac{1}{R} \sum_{j=1}^s \sum_{R_0 < |n| < R} |\widehat{u}_j(n)|^2 = 0,$$

there exists a constant  $R_1 > R_0$  such that  $\widehat{u}(n) = 0$  for  $|n| > R_1$  i.e.  $\# \text{supp} \widehat{u} < \infty$ .

## Application for eigenvalue problems (UCP of Helmholtz type eq.)

### Assumption (unique continuation property on lattices)

(A-4) If  $\hat{u}$  satisfies

$$(\widehat{H}_0 + \widehat{V} - \lambda)\hat{u} = 0 \quad \text{on } \mathcal{V}_0$$

and, for a constant  $R_1 > 0$ ,  $\hat{u}(n) = 0$ ,  $|n| > R_1$ , then  $\hat{u} = 0$  on whole  $\mathcal{V}_0$ .

- ▶ UCP on lattices is slightly different from elliptic PDE on  $\mathbf{R}^d$  or manifolds.
- ▶ If we assume (A-1),  $\lambda \notin \mathcal{T}_1$  and  $\widehat{V} = 0$ , UCP holds on our examples.
- ▶ If  $\widehat{V} \neq 0$ , it is **not** sufficient to assume (A-1) and  $\lambda \notin \mathcal{T}_1$ . In fact, on kagome lattice and subdivision lattice, UCP does not holds for any  $\lambda \in \mathbf{R}$ . Moreover, for any  $\lambda \in \mathbf{R}$ , we can construct  $\widehat{V}$  such that  $\lambda \in \sigma_p(\widehat{H}_0 + \widehat{V})$ .

## Absence of embedded eigenvalues

### Theorem

We assume (A-1) and (A-4).

If  $\lambda \in \sigma_{ess}(\widehat{H}) \setminus \mathcal{T}_1$ , we have  $\lambda \notin \sigma_p(\widehat{H})$  i.e.

$$\sigma_p(\widehat{H}) \cap (\sigma_{ess}(\widehat{H}) \setminus \mathcal{T}_1) = \emptyset.$$

Sketch of proof.

- ▶ Since  $\sharp \text{supp} \widehat{V} < \infty$ , we can apply the Rellich type theorem for the eigenfunction  $\widehat{\psi}_\lambda \in \ell^2(\mathcal{V}_0)$ .
- ▶ For a sufficiently large  $R_1 > 0$ ,  $\widehat{\psi}_\lambda(n) = 0$  for  $|n| > R_1$ .
- ▶ From (A-4),  $\widehat{\psi}_\lambda(n) = 0$  on  $\mathcal{V}_0$ . This is a contradiction.

## Examples of $\mathcal{T}_1$ (1)

(1) : Square lattice

- ▶  $\sigma(\widehat{H}_0) = [-1, 1]$
- ▶  $\mathcal{T}_1 = \{-1, 1\}$

(2) : Hexagonal lattice

- ▶  $\sigma(\widehat{H}_0) = [-1, 1]$
- ▶  $\mathcal{T}_1 = \{-1, 0, 1\}$

(3) : Kagome lattice

- ▶  $\sigma(\widehat{H}_0) = [-1, 1/2]$
- ▶  $\mathcal{T}_1 = \{-1, -1/4, 1/2\}, 1/2 \in \sigma_p(\widehat{H}_0).$

(4) : subdivision

- ▶  $\sigma(\widehat{H}_0) = [-1, 1]$
- ▶  $\mathcal{T}_1 = \{-1, 0, 1\}, 0 \in \sigma_p(\widehat{H}_0).$

## Examples of $\mathcal{T}_1$ (2)

(5) : Ladder

- ▶  $\sigma(\widehat{H}_0) = [-1, 1]$
- ▶  $\mathcal{T}_1 = \left\{ \frac{2d-1}{2d+1} \leq |\lambda| \leq 1 \right\}$

(6) : Graphite

- ▶  $\sigma(\widehat{H}_0) = [-1, 1]$
- ▶  $\mathcal{T}_1 = \{1/2 \leq |\lambda| \leq 1\}$



## Procedure of UCP

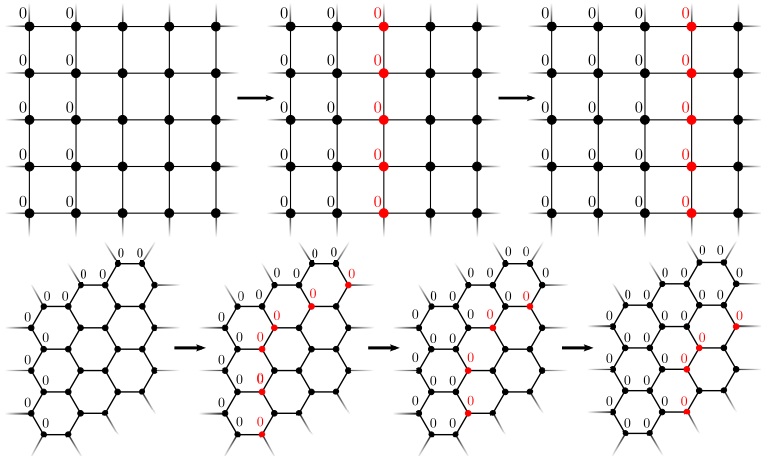


Figure: Unique continuation on lattices

## Example of embedded eigenvalue -Kagome lattice-

If we put  $\widehat{V}(v) = \alpha$  for  $v = x_1, \dots, x_6$ , else  $\widehat{V}(v) = 0$ , an eigenfunction satisfies  $\widehat{u}(v) = (-1)^j$  for  $v = x_1, \dots, x_6$ , else  $\widehat{u}(v) = 0$  with the eigenvalue  $\lambda = \alpha + 1/2$ .

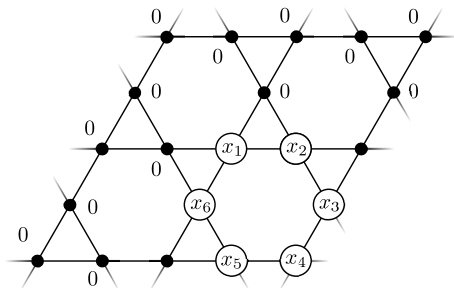


Figure: Eigenfunction with an embedded eigenvalue

## Example of embedded eigenvalue -Subdivision lattice-

If we put  $\widehat{V}(v) = \alpha$  for  $v = x_1, \dots, x_4$ , else  $\widehat{V}(v) = 0$ , an eigenfunction satisfies  $\widehat{u}(v) = (-1)^j$  for  $v = x_1, \dots, x_4$ , else  $\widehat{u}(v) = 0$  with the eigenvalue  $\lambda = \alpha$ .

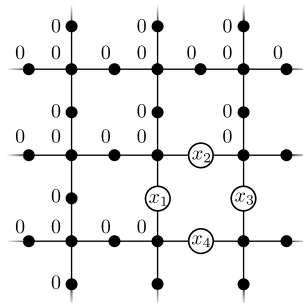


Figure: Eigenfunction with an embedded eigenvalue

## Example of embedded eigenvalue -ladder-

For any constant  $\alpha \neq 0$ , we put  $\hat{V}(0) = \alpha I_2$  and  $\hat{V}(n) = 0$  for  $n \neq 0$ .  
 For any  $\lambda \in \mathcal{T}_1$ , we can choose  $\alpha$  and construct an eigenfunction which decays at infinity.

Graphite is similar.

## Sketch of proof for Rellich type theorem (1)

- ▶ Multiplying the co-factor of  $H_0(x) - \lambda$ , we can diagonalize the equation as

$$p(x, \lambda)I_s u = g,$$

so that we pick up a component. In the following, we deal it with a single equation.

- ▶ We obtain  $u \in C^\infty(\mathbb{T}^d \setminus M_{\lambda, \text{sing}}^C)$ . In particular, we have  $g(x) = 0$  on  $M_{\lambda, \text{reg}}^C \cap \mathbb{T}^d$ .
- ▶ In view of (A-1-2), we can extend analytically  $g(z) = 0$  to  $M_{\lambda, \text{reg}}^C$ .
- ▶ Hence  $g(z)/p(z, \lambda)$  is analytic in a neighborhood of  $M_{\lambda, \text{reg}}^C$ .
- ▶ From (A-1-1),  $M_{\lambda, \text{sing}}^C$  is a removable singularity, so that  $g(z)/p(z, \lambda)$  is analytic in  $\mathbb{T}_C^d$ .

## Sketch of proof for Rellich type theorem (2)

- ▶ Changing the variable  $w_j = e^{iz_j}$ , we have

$$\frac{g(z)}{p(z, \lambda)} = \frac{G(w)}{P(w, \lambda)} \prod_{j=1}^d w_j^{\gamma_j - \beta_j}, \quad G, P \in C[w_1, \dots, w_d].$$

- ▶ Since LHS is analytic,  $G/P$  is also analytic. In particular,  $G(w) = 0$  on  $\{w \in \mathbb{C}^d ; P(w, \lambda) = 0\}$ .
- ▶ Hilbert Nullstellensatz implies that  $P$  divides  $G$ .
- ▶ Therefore  $u = g/p$  is a trigonometric polynomial. This implies that  $\hat{u} = \mathcal{U}_{\mathcal{L}_0}^* u$  has a finite support.