Quantum White Noise Derivatives and Implementation Problem

On the occasion of thier 60th birthdays of Professors K. R. Ito and I. Ojima

Nobuaki Obata

GSIS, Tohoku University

Fukuoka, November 29, 2009

The Implementation Problem

 $a(\xi), a^*(\eta)$: annihilation and creation operators on Boson Fock space $\Gamma(H)$ satisfying

$$\mathsf{CCR}: \quad [a(\xi),a(\eta)] = [a^*(\xi),a^*(\eta)] = 0, \quad [a(\xi),a^*(\eta)] = \langle \xi,\eta\rangle$$

Consider transformed annihilation and creation operators:

$$b(\zeta)=a(S\zeta)+a^*(T\zeta), \hspace{1em} b^*(\zeta)=a^*(S\zeta)+a(T\zeta).$$

The implementation problem [Berezin (1966), Ruijsenaars (1977), ...] is to find a (unitary) operator U on the Boson Fock space $\Gamma(H)$ satisfying

$$\begin{array}{ccc} \Gamma(H) & \stackrel{U}{\longrightarrow} & \Gamma(H) & & \Gamma(H) & \stackrel{U}{\longrightarrow} & \Gamma(H) \\ \hline a(\zeta) & & & \downarrow^{b(\zeta)} & & a^*(\zeta) & & \downarrow^{b^*(\zeta)} \\ \Gamma(H) & \stackrel{U}{\longrightarrow} & \Gamma(H) & & & \Gamma(H) & \stackrel{U}{\longrightarrow} & \Gamma(H) \end{array}$$

 $\underline{\text{Remarks:}} (1) [b(\zeta), b(\eta)] = [b^*(\zeta), b^*(\eta)] = 0 \iff T^*S = S^*T$ $(2) [b(\zeta), b^*(\eta)] = \langle \zeta, \eta \rangle \iff S^*S - T^*T = I$

Plan

[0]

1. Quantum White Noise Calculus

- 1.1. Background and Notation
- 1.2. White Noise Operators
- 1.3. Quantum White Noise
- 1.4. Integral Kernel Operators and Fock Expansion
- 2. Quantum White Noise Derivatives
 - 2.1. Definition
 - 2.2. Examples
 - 2.3. Wick Product
 - 2.4. Wick Derivations
- 3. Differential Equations for White Noise Operators
 - 3.1. Differential Equations
 - 3.2. Reproducing Irreducubility of CCR
 - 3.3. Linear Equations
- 4. Implementation Problem for CCR
 - 4.1. The Implementation Problem
 - 4.2. Our Approach
 - 4.3. Solution to the Implementation Problem

1. Quantum White Noise Calculus

1.1. Background and Notation

The Boson Fock space over $H = L^2(T)$ is defined by

$$\Gamma(H)=\left\{\phi=(f_n)\,;\,f_n\in H^{\widehat{\otimes}n}\,,\,\,\|\phi\|^2=\sum_{n=0}^\infty n!|f_n|_0^2<\infty
ight\},$$

where T is a topological space equipped with a σ -finite Borel measure dt, $|f_n|_0$ is the usual L^2 -norm of $H^{\widehat{\otimes}n} = L^2_{\mathrm{sym}}(T^n)$.

The annihilation and creation operator at a point $t \in T$

$$egin{aligned} a_t:(0,\ldots,0,m{\xi}^{\otimes n},0,\ldots)&\mapsto(0,\ldots,0,nm{\xi}(t)m{\xi}^{\otimes(n-1)},0,0,\ldots)\ a_t^*:(0,\ldots,0,m{\xi}^{\otimes n},0,\ldots)&\mapsto(0,\ldots,0,0,m{\xi}^{\otimes n}\widehat{\otimes}\delta_t,0,\ldots) \end{aligned}$$

A "general" Fock space operator takes the form:

$$\sum_{l,m=0}^{\infty}\int_{T^{l+m}} \kappa_{l,m}(s_1,\ldots,s_l,t_1,\ldots,t_m) a_{s_1}^*\cdots a_{s_l}^*a_{t_1}\cdots a_{t_m} ds_1\cdots ds_l dt_1\cdots dt_m$$

Quantum field theory: e.g., Haag (1955), Berezin (1966), Krée (1988), etc.

1.2. White Noise Operators

I) Gelfand triple for $H = L^2(T)$:

$$E\subset H=L^2(T)\subset E^*,\qquad E=\operatorname*{proj}_{p o\infty}E_p\,,\quad E^*=\operatorname*{ind}_{p o\infty}E_{-p}\,,$$

where E_p is a dense subspace of H and is a Hilbert space for itself.

II) Gelfand triple for $\Gamma(H)$ (e.g., Hida–Kubo–Takenaka space (1980)):

$$(E) \subset \Gamma(H) \subset (E)^*, \qquad (E) = \operatorname{proj}_{p \to \infty} \lim \Gamma(E_p), \quad (E)^* = \operatorname{ind}_{p \to \infty} \lim \Gamma(E_{-p}),$$

Note: (1) $\Gamma(H) \cong L^2(E^*, \mu)$ (Wiener–Itô–Segal isomorphism) (2) (E) is the space of test functions and (E)^{*} the space of distributions.

Definition (White noise operator)

A continuous operator from (E) into $(E)^*$ is called a <u>white noise operator</u>. Let $\mathcal{L}((E), (E)^*)$ denote the space of white noise operators, equipped with the topology of bounded convergence.

Note: $\mathcal{L}((E), (E))$, $\mathcal{L}((E)^*, (E)^*)$ and $\mathcal{B}(\Gamma(H))$ are subspaces of $\mathcal{L}((E), (E)^*)$.

Theorem (Quantum white noise is very regular)

 $a_t \in \mathcal{L}((E), (E))$ and $a_t^* \in \mathcal{L}((E)^*, (E)^*)$ for all $t \in \mathbb{R}$. Moreover, both maps $t \mapsto a_t \in \mathcal{L}((E), (E))$ and $t \mapsto a_t^* \in \mathcal{L}((E)^*, (E)^*)$ are operator-valued rapidly decreasing functions, i.e., belongs to $E \otimes \mathcal{L}((E), (E))$ and $E \otimes \mathcal{L}((E)^*, (E)^*)$, respectively. (The pair $\{a_t, a_t^*; t \in T\}$ is called the quantum white noise on T.)

Smeared operators

$$a(\zeta) = \int \zeta(t) a_t \, dt, \qquad a^*(\zeta) = \int \zeta(t) a_t^* \, dt$$

Traditional approach

- ζ is a test function, e.g., $\zeta \in \mathcal{S}(\mathbb{R})$.
- $a(\zeta), a^*(\zeta)$ are unbounded opereators in $\Gamma(H)$.

White noise approach

- ζ is a distribution, e.g., $\zeta \in \mathcal{S}'(\mathbb{R})$.
- **2** $a(\zeta)$, $a^*(\zeta)$ are white noise operators, i.e., belong to $\mathcal{L}((E), (E)^*)$.
- ${f 0}$ In fact, $a(\zeta)\in {\cal L}((E),(E))$ and $a^*(\zeta)\in {\cal L}((E)^*,(E)^*).$

Definition (Integral kernel operator)

Given $\kappa_{l,m} \in (E^{\otimes (l+m)})^*$, $l,m=0,1,2,\ldots$, the integral kernel operator

$$egin{aligned} \Xi_{l,m}(\kappa_{l,m}) \ &= \int_{T^{l+m}} & \kappa_{l,m}(s_1,\cdots,s_l,t_1,\cdots,t_m) a^*_{s_1}\cdots a^*_{s_l}a_{t_1}\cdots a_{t_m}ds_1\cdots ds_l dt_1\cdots dt_m \end{aligned}$$

is defined and is a white noise operator, i.e., $\Xi_{l,m}(\kappa_{l,m})\in\mathcal{L}((E),(E)^*).$

Theorem (O.(1993), cf. Berezin (1966), Krée (1988))

Every white noise operator $\Xi \in \mathcal{L}((E), (E)^*)$ admits the infinite series expansion:

$$\Xi = \sum_{l,m=0}^\infty \Xi_{l,m}(\kappa_{l,m}), \qquad \kappa_{l,m} \in (E^{\otimes (l+m)})^*,$$

where the right-hand side converges in $\mathcal{L}((E), (E)^*)$. If $\Xi \in \mathcal{L}((E), (E))$, then $\kappa_{l,m} \in E^{\otimes l} \otimes (E^{\otimes m})^*$ and the series converges in $\mathcal{L}((E), (E))$.

2. Quantum White Noise Derivatives

2.1. Definition

For a Brownian (or white noise) function Φ stochastic derivatives (gradients) were introduced by Malliavin, Hida, Gross, ...

۲

$$abla \Phi, \quad rac{\delta \Phi}{\delta \dot{B}(t)}\,, \quad \partial_t \Phi, \quad a_t \Phi$$

A quantum counterpart

A white noise operator Ξ is considered as a function of quantum white noise: $\Xi = \Xi(a_s, a_t^*; s, t \in T)$. We should like to define the derivatives with respect to a_s and a_t^* :

$$rac{\delta \Xi}{\delta a_s}$$
 and $rac{\delta \Xi}{\delta a_t^*}$

Expected properties:

$$egin{aligned} &rac{\delta}{\delta a_s}\int f(t)a_tdt = f(s)I\ &rac{\delta}{\delta a_s}\int f(s,t)a_sa_tdsdt = \int f(s,t)a_t\,dt + \int f(t,s)a_t\,dt\ &rac{\delta}{\delta a_t^*}\int f(s,t)a_sa_t^*dsdt = \int f(s,t)a_s\,ds \end{aligned}$$

2.1. Definition

Definition (Ji–Obata (2007))

For $\Xi \in \mathcal{L}((E), (E)^*)$ and $\zeta \in E$ we define $D_{\zeta}^{\pm}\Xi \in \mathcal{L}((E), (E)^*)$ by

$$D^+_\zeta\Xi=[a(\zeta),\Xi], \qquad D^-_\zeta\Xi=-[a^*(\zeta),\Xi].$$

These are called the <u>creation derivative</u> and <u>annihilation derivative</u> of Ξ , respectively. Both together are called the *quantum white noise derivatives*.

<u>Note</u>: For $\zeta \in E$, both

$$a(\zeta) = \Xi_{0,1}(\zeta) = \int_T \zeta(t) a_t \, dt, \quad a^*(\zeta) = \Xi_{1,0}(\zeta) = \int_T \zeta(t) a_t^* \, dt,$$

belong to $\mathcal{L}((E),(E))\cap\mathcal{L}((E)^*,(E)^*).$

- $(D_{\zeta}^+\Xi)^* = D_{\zeta}^-(\Xi^*)$ and $(D_{\zeta}^-\Xi)^* = D_{\zeta}^+(\Xi^*)$.
- **2** D_{ζ}^{\pm} is a continuous linear map from $\mathcal{L}((E), (E)^*)$ into itself.
- Moreover, (ζ,Ξ) → D[±]_ζΞ is a continuous bilinear map from E × L((E), (E)*) into L((E), (E)*).

Remark: Pointwisely Defined QWN-Derivatives

Recall: The smeared annihilation and creation operators

$$a(f)=\int_T f(t)a_t\,dt,\qquad a^*(f)=\int_T f(t)a_t^*\,dt\,.$$

It is natural to introduce D_t^{\pm} to have

$$D_{\zeta}^+ = \int_T \zeta(t) D_t^+ dt, \qquad D_{\zeta}^- = \int_T \zeta(t) D_t^- dt.$$

In fact, this expression is useful for computation.

② However, it is not straightforward to define D_t^\pm for each point $t\in T$ because

$$D_t^+\Xi=[a_t,\Xi]=a_t\Xi-\Xi a_t, \qquad D_t^-\Xi=-[a_t^*,\Xi]=-a_t^*\Xi+\Xi a_t^*$$

are not well-defined in general.

Nevertheless, the pointwisely defined quantum white noise derivatives D[±]_t are well formulated for <u>admissible white noise operators</u> L(G, G*) [Ji-Obata, 2009, to appear].

2.2. Examples

The canonical correspondence (kernel theorem) between $S \in \mathcal{L}(E, E^*)$ and $\tau = \tau_S \in (E \otimes E)^*$ is given by $\langle \tau_S, \eta \otimes \xi \rangle = \langle S\xi, \eta \rangle$ for $\xi, \eta \in E$.

(1) The generalized Gross Laplacian associated with S is defined by

$$\Delta_{\mathrm{G}}(S) = \Xi_{0,2}(au_S) = \int_{T imes T} au_S(s,t) a_s a_t \, ds dt$$

Note that $\Delta_{\mathbf{G}}(S) \in \mathcal{L}((E), (E))$. Then,

$$D^+_\zeta \Delta_{\mathrm{G}}(S) = 0, \qquad D^-_\zeta \Delta_{\mathrm{G}}(S) = a(S\zeta) + a(S^*\zeta)$$

In fact, since

$$D^-_t\Delta_{\mathrm{G}}(S) = \int_T au_S(s,t) a_s\,ds + \int_T au_S(t,s) a_s\,ds$$

we have

$$egin{aligned} D_{\zeta}^{-}\Delta_{\mathrm{G}}(S) &= \int_{T imes T} au_{S}(s,t)a_{s}\zeta(t)\,dsdt + \int_{T imes T} au_{S}(t,s)a_{s}\zeta(t)\,dsdt \ &= \int_{T}S\zeta(s)a_{s}\,ds + \int_{T}S^{*}\zeta(s)a_{s}\,ds = a(S\zeta) + a(S^{*}\zeta) \end{aligned}$$

2.2. Examples

(2) The adjoint of $\Delta_{\mathrm{G}}(S) \in \mathcal{L}((E)^*, (E)^*)$ is given by

$$\Delta^*_{\mathrm{G}}(S) = \Xi_{2,0}(au_S) = \int_{T imes T} au_S(s,t) a^*_s a^*_t \, ds dt$$

The quantum white noise derivatives are given by

$$D^-_\zeta\Delta^*_\mathrm{G}(S)=0, \qquad D^+_\zeta\Delta^*_\mathrm{G}(S)=a^*(S\zeta)+a^*(S^*\zeta)$$

(3) The conservation operator associated with S is defined by

$$\Lambda(S) = \Xi_{1,1}(au_S) = \int_{T imes T} au_S(s,t) a_s^* a_t \, ds dt$$

In general, $\Lambda(S) \in \mathcal{L}((E), (E)^*)$.

The quantum white noise derivatives are given by

$$D_{\zeta}^{-}\Lambda(S) = a^{*}(S\zeta), \qquad D_{\zeta}^{+}\Lambda(S) = a(S^{*}\zeta).$$

2.3. Wick Product

The <u>Wick product</u> of white noise operators $\Xi_1, \Xi_2 \in \mathcal{L}((E), (E)^*)$, denoted by $\Xi_1 \diamond \Xi_2$, is characterized by

$$a_t \diamond \Xi = \Xi \diamond a_t = \Xi a_t\,, \qquad a_t^* \diamond \Xi = \Xi \diamond a_t^* = a_t^* \Xi$$

Equipped with the Wick product, $\mathcal{L}((E), (E)^*)$ becomes a commutative algebra.

Definition (Wick product)

For $\Xi_1, \Xi_2 \in \mathcal{L}((E), (E)^*)$ the Wick (or normal-ordered) product $\Xi_1 \diamond \Xi_2$ is defined by

$$(\Xi_1\diamond\Xi_2)^{\widehat{}}(\xi,\eta)=\widehat{\Xi}_1(\xi,\eta)\widehat{\Xi}_2(\xi,\eta)e^{-\langle\xi,\eta
angle},\qquad \xi,\eta\in E$$

where $\widehat{\Xi}(\xi,\eta)$ is the symbol of a white noise operator $\Xi\in\mathcal{L}((E),(E)^*)$ defined by

$$\widehat{\Xi}(\xi,\eta) = \langle\!\langle \Xi \phi_{\xi}, \phi_{\eta}
angle\!
angle, \qquad \xi,\eta \in E,$$

where $\phi_{\xi} = (1, \xi, \dots, \xi^{\otimes n}/n!, \dots)$ is an *exponential vector*. This is verified by the characterization theorem for operator symbols (see O. LNM 1577 (1994)]

15 / 35

2.4. Wick Derivations

 $(\mathcal{L}((E),(E)^*),\diamond)$ is a commutative algebra.

Definition (Wick derivation)

A continuous linear map $\mathcal{D}: \mathcal{L}((E), (E)^*) \to \mathcal{L}((E), (E)^*)$ is called a <u>Wick derivation</u> if

$$\mathcal{D}(\Xi_1 \diamond \Xi_2) = (\mathcal{D}\Xi_1) \diamond \Xi_2 + \Xi_1 \diamond (\mathcal{D}\Xi_2)$$

for all $\Xi_1, \Xi_2 \in \mathcal{L}((E), (E)^*)$.

Theorem

The creation and annihilation derivatives D_{ζ}^{\pm} are Wick derivations for any $\zeta \in E$.

Note: It is proved that a general Wick derivation \mathcal{D} is expressed in the form:

$$\mathcal{D} = \int_T F(t) \diamond D_t^+ dt + \int_T G(t) \diamond D_t^- dt,$$

where $F,G\in E\otimes \mathcal{L}((E),(E)^*).$

Proof.

In general, for $\Xi \in \mathcal{L}((E),(E)^*)$ we have

$$(D_{\zeta}^{+}\Xi)^{\widehat{}}(\xi,\eta) = \langle\!\langle (a(\zeta)\Xi - \Xi a(\zeta))\phi_{\xi},\phi_{\eta}\rangle\!\rangle$$

$$= \langle\!\langle \Xi\phi_{\xi},a^{*}(\zeta)\phi_{\eta}\rangle\!\rangle - \langle\!\langle \Xi a(\zeta)\phi_{\xi},\phi_{\eta}\rangle\!\rangle$$

$$= \left.\frac{d}{dt}\right|_{t=0} \langle\!\langle \Xi\phi_{\xi},\phi_{\eta+t\zeta}\rangle\!\rangle - \langle\xi,\zeta\rangle\langle\!\langle \Xi\phi_{\xi},\phi_{\eta}\rangle\!\rangle$$

$$= \left.\frac{d}{dt}\right|_{t=0} \widehat{\Xi}(\xi,\eta+t\zeta) - \langle\xi,\zeta\rangle\widehat{\Xi}(\xi,\eta).$$
(1)

Then for $\Xi=\Xi_1\diamond\Xi_2$ we have

$$\begin{split} (D_{\zeta}^{+}\Xi)^{\widehat{}}(\xi,\eta) &= \left. \frac{d}{dt} \right|_{t=0} \widehat{\Xi}_{1}(\xi,t\zeta+\eta) \widehat{\Xi}_{2}(\xi,t\zeta+\eta) e^{-\langle\xi,t\zeta+\eta\rangle} \\ &- \langle\xi,\zeta\rangle \widehat{\Xi}_{1}(\xi,\eta) \widehat{\Xi}_{2}(\xi,\eta) e^{-\langle\xi,\eta\rangle} \\ &= \left(\left. \frac{d}{dt} \right|_{t=0} \widehat{\Xi}_{1}(\xi,t\zeta+\eta) \right) \widehat{\Xi}_{2}(\xi,\eta) e^{-\langle\xi,\eta\rangle} \\ &+ \widehat{\Xi}_{1}(\xi,\eta) \left(\left. \frac{d}{dt} \right|_{t=0} \widehat{\Xi}_{2}(\xi,t\zeta+\eta) \right) e^{-\langle\xi,\eta\rangle} \\ &- 2\langle\xi,\zeta\rangle \widehat{\Xi}_{1}(\xi,\eta) \widehat{\Xi}_{2}(\xi,\eta) e^{-\langle\xi,\eta\rangle}. \end{split}$$

Viewing (1) once again, we obtain

$$(D_{\zeta}^+\Xi)^{\hat{}}(\xi,\eta) = ((D_{\zeta}^+\Xi_1)\diamond\Xi_2)^{\hat{}}(\xi,\eta) + (\Xi_1\diamond(D_{\zeta}^+\Xi_2))^{\hat{}}(\xi,\eta).$$

3. Differential Equations for White Noise Operators

A general form

 $\mathcal{D}: \mathcal{L}((E), (E)^*)
ightarrow \mathcal{L}((E), (E)^*)$: a Wick derivation $f: \mathcal{L}((E), (E)^*)
ightarrow \mathcal{L}((E), (E)^*)$: a map

 $\mathcal{D}\Xi=f(\Xi)$

Simple cases:

- **1** $\mathcal{D}\Xi = 0$ ("constant" with respect to \mathcal{D})
- **2** $\mathcal{D}\Xi = G \diamond \Xi$ with $G \in \mathcal{L}((E), (E)^*)$ (linear equation)

General cases: interesting for characterizing white noise operators (future problem)?

3.2. Reproducing Irreducubility of CCR

Let us consider the (system of) differential equations:

$$D_{\zeta}^{+}\Xi = 0, \qquad \zeta \in E.$$

We expect easily that $\Xi = \Xi(a_s, a_t^*; s, t \in T)$ does not depend on the creation operators. In fact, by Fock expansion we see that the solutions to (2) are given by

$$\Xi = \sum_{m=0}^\infty \Xi_{0,m}(\kappa_{0,m}).$$

In a similar manner, the solutions to

$$D_{\zeta}^{-}\Xi=0,\qquad \zeta\in E,$$

are given by

$$\Xi = \sum_{l=0}^{\infty} \Xi_{l,0}(\kappa_{l,0}).$$

Consequently, a white noise operator satisfying both (2) and (3) are the scalar operators. Thus, the irreducibility of the canonical commutation relation is reproduced.

3.2. Linear Equations

Given a Wick derivation \mathcal{D} and $G \in \mathcal{L}((E), (E)^*)$, consider

$$\mathcal{D}\Xi = G \diamond \Xi$$

The Wick exponential is defined by

$$ext{wexp } Y = \sum_{n=0}^{\infty} \frac{1}{n!} \, Y^{\diamond n}, \qquad Y \in \mathcal{L}((E), (E)^*),$$

whenever the series converges in $\mathcal{L}((E), (E)^*)$.

Theorem

Every solution to (4) is of the form:

$$\Xi = (\text{wexp } Y) \diamond F, \tag{5}$$

where (i)
$$Y \in \mathcal{L}((E), (E)^*)$$
 is a solution to $\mathcal{D}Y = G$;
(ii) wexp Y should be defined in $\mathcal{L}((E), (E)^*)$;
(iii) $F \in \mathcal{L}((E), (E)^*)$ is arbitrary satisfying $\mathcal{D}F = 0$.

(4)

Proof.

It is straightforward to see that

$$\Xi = (\mathrm{wexp} \ Y) \diamond F$$

is a solution to

$$\mathcal{D}\Xi = G \diamond \Xi \tag{6}$$

To prove the converse, let Ξ be an arbitrary solution to (6). Set

$$F = (\operatorname{wexp}(-Y)) \diamond \Xi$$

Obviously, $F \in \mathcal{L}((E), (E)^*)$ and $\Xi = (\text{wexp } Y) \diamond F$. We only need to show that $\mathcal{D}F = 0$. In fact,

$$\mathcal{D}F = -\mathcal{D}Y \diamond (\operatorname{wexp}(-Y)) \diamond \Xi + (\operatorname{wexp}(-Y)) \diamond \mathcal{D}\Xi$$

= $-G \diamond (\operatorname{wexp}(-Y)) \diamond \Xi + (\operatorname{wexp}(-Y)) \diamond G \diamond \Xi = 0.$

This completes the proof.

Example (1)

$$D_\zeta^-\Xi=2a(\zeta)\diamond\Xi,\qquad \zeta\in E.$$

 $\textbf{ let we need to find } Y \in \mathcal{L}((E),(E)^*) \text{ satisfying } D_{\zeta}^-Y = 2a(\zeta).$

In fact,

$$Y = \Delta_G = \int a_t^2 \, dt$$

is a solution.

- **(a)** Moreover, it is easily verified that wexp $\Delta_{\mathbf{G}}$ is defined in $\mathcal{L}((E), (E))$.
- Then, a general solution to (7) is of the form:

$$\Xi = (\text{wexp } \Delta_{\mathbf{G}}) \diamond F, \tag{8}$$

where $D_{\zeta}^{-}F = 0$ for all $\zeta \in E$.

Example (2)

$$egin{aligned} D_{\zeta}^- \Xi &= 2a(\zeta) \diamond \Xi, \qquad \zeta \in E, \ D_{\zeta}^+ \Xi &= 0. \end{aligned}$$

(9)

By Example (1) the solution is of the form:

$$\Xi = (ext{wexp} \ \Delta_{ ext{G}}) \diamond F, \qquad D_{\zeta}^{-}F = 0 ext{ for all } \zeta \in E.$$

2 We need only to find additional conditions for F satisfying $D_{\zeta}^{+}\Xi = 0$.

 ${f 0}$ Noting that $D_\zeta^+\Delta_{
m G}=0$, we have

$$D^+_\zeta \Xi = (ext{wexp } \Delta_{\mathrm{G}}) \diamond D^+_\zeta F = 0.$$

Hence $D_{\zeta}^+ F = 0$ for all $\zeta \in E$. Consequently, F is a scalar operator (irreducubility of CCR).

Finally, the solution to (9) is of the form:

$$\Xi = C ext{ wexp } \Delta_{\mathrm{G}}, \qquad C \in \mathbb{C}.$$

4. Implementation Problem for CCR

4.1. The Implementation Problem

Let $S, T \in \mathcal{L}(E, E)$ and consider transformed annihilation and creation operators:

$$b(\zeta)=a(S\zeta)+a^*(T\zeta), \ \ b^*(\zeta)=a^*(S\zeta)+a(T\zeta),$$

where $\zeta \in E$. We know that $b(\zeta), b^*(\zeta) \in \mathcal{L}((E), (E)) \cap \mathcal{L}((E)^*, (E)^*)$.

The implementation problem

is to find a white noise operator $U \in \mathcal{L}((E),(E)^*)$ satisfying

$$(E) \xrightarrow{U} (E)^{*} \qquad (E) \xrightarrow{U} (E)^{*}$$
$$a(\zeta) \downarrow \qquad \qquad \downarrow b(\zeta) \qquad \qquad a^{*}(\zeta) \downarrow \qquad \qquad \downarrow b^{*}(\zeta)$$
$$(E) \xrightarrow{U} (E)^{*} \qquad (E) \xrightarrow{U} (E)^{*}$$

<u>Remarks</u>: (1) $T^*S = S^*T$ is equivalent to

$$[b(\zeta),b(\eta)]=[b^*(\zeta),b^*(\eta)]=0,\qquad \zeta,\eta\in E.$$

(2) $S^*S - T^*T = I$ is equivalent to

$$[b(\zeta),b^*(\eta)]=\langle \zeta,\eta
angle,\qquad \zeta,\eta\in E.$$

4.2. Our Approach

$$egin{aligned} Ua(\zeta) &= b(\zeta)U \ &= (a(S\zeta) + a^*(T\zeta))\,U \ &= D^+_{S\zeta}U + Ua(S\zeta) + a^*(T\zeta)U, \ D^+_{S\zeta}U &= Ua(\zeta) - Ua(S\zeta) - a^*(T\zeta)U \ &= Ua(\zeta - S\zeta) - a^*(T\zeta)U \ &= [a(\zeta - S\zeta) - a^*(T\zeta)] \diamond U. \end{aligned}$$

Thus,

$$Ua(\zeta)=b(\zeta)U \quad \Longleftrightarrow \quad D^+_{S\zeta}U=[a(\zeta-S\zeta)-a^*(T\zeta)]\diamond U.$$

Similarly,

$$Ua^*(\zeta) = b^*(\zeta)U \quad \Longleftrightarrow \quad (D_\zeta^- - D_{T\zeta}^+)U = [a^*(S\zeta - \zeta) + a(T\zeta)] \diamond U.$$

Theorem

Assume that S is invertible and that $T^*S = S^*T$. Then a white noise operator $U \in \mathcal{L}((E), (E)^*)$ satisfies the intertwining property:

$$Ua(\zeta)=b(\zeta)U,\qquad \zeta\in E,$$

if and only if U is of the form

$$U = \operatorname{wexp}\left\{-\frac{1}{2}\Delta_{\mathrm{G}}^{*}(TS^{-1}) + \Lambda((S^{-1})^{*} - I)\right\} \diamond F,$$
 (10)

where $F \in \mathcal{L}((E), (E)^*)$ fulfills $D_{\zeta}^+ F = 0$ for all $\zeta \in E$.

Remark:

$$\begin{split} & \text{wexp } \left\{ -\frac{1}{2} \Delta_{\text{G}}^{*}(TS^{-1}) \right\} = e^{-\frac{1}{2} \Delta_{\text{G}}^{*}(TS^{-1})} \\ & \text{wexp } \left\{ \Lambda((S^{-1})^{*} - I) \right\} = \Gamma((S^{-1})^{*}) \\ & U = e^{-\frac{1}{2} \Delta_{\text{G}}^{*}(TS^{-1})} \Gamma((S^{-1})^{*}) F \quad \text{(usual composition)} \end{split}$$

Proof.

We only need to solve the differential equation

$$D_{S\zeta}^+ U = [a(\zeta - S\zeta) - a^*(T\zeta)] \diamond U.$$
(11)

We readily know that

$$D^+_{S\zeta}\Lambda((S^{-1})^*-I) = a(\zeta-S\zeta), \qquad D^+_{S\zeta}\Delta^*_{
m G}(TS^{-1}) = 2a^*(T\zeta).$$

Then by the general result a general form of the solutions to (11) is given by

$$U = \operatorname{wexp} \left\{ -\frac{1}{2} \Delta_{\mathrm{G}}^{*}(TS^{-1}) + \Lambda((S^{-1})^{*} - I) \right\} \diamond F,$$

where $F \in \mathcal{L}((E), (E)^*)$ is an arbitrary white noise operator satisfying $D^+_{S\zeta}F = 0$ for all $\zeta \in E$.

• Since S is invertible, the last condition for F is equivalent to that $D_{\zeta}^+F=0$ for all $\zeta\in E.$

Theorem

Assume the following conditions:

- (i) S is invertible;
- (ii) $T^*S = S^*T;$
- (iii) $S^*S T^*T = I;$
- (iv) $ST^* = TS^*$.

Then a white noise operator $U \in \mathcal{L}((E), (E)^*)$ satisfies the intertwining property:

$$Ua^*(\zeta)=b^*(\zeta)U,\qquad \zeta\in E,$$

if and only if U is of the form:

$$U = ext{wexp} \, \left\{ -rac{1}{2} \Delta_{\mathrm{G}}^{*}(TS^{-1}) + \Lambda((S^{-1})^{*} - I) + rac{1}{2} \Delta_{\mathrm{G}}(S^{-1}T)
ight\} \diamond G,$$

where $G \in \mathcal{L}((E),(E)^*)$ is an arbitrary white noise operator satisfying

$$(D_\zeta^- - D_{T\zeta}^+)G = 0 \quad \text{for all } \zeta \in E.$$

Proof.

Our task is to solve the differential equation:

$$(D_\zeta^- - D_{T\zeta}^+)U = [a^*(S\zeta - \zeta) + a(T\zeta)] \diamond U.$$

I First we need to find a solution to the differential equation:

$$(D_{\zeta}^{-} - D_{T\zeta}^{+})Y = a^{*}(S\zeta - \zeta) + a(T\zeta).$$
(12)

As is easily verified,

$$Y=\Delta^*_{
m G}(K)+\Lambda(L)+\Delta_{
m G}(M),\qquad K=K^*,\quad M=M^*,$$

satisfies (12) if and only if

$$2M - L^*T = T, \qquad L - 2KT = S - I.$$

Thanks to the conditions (i)–(iv) we may choose

$$K = -rac{1}{2} TS^{-1}, \quad L = (S^{-1})^* - I, \quad M = rac{1}{2} S^{-1}T.$$

In the assertion follows immediately from our general theorem.

Theorem

Assume the following conditions:

- (i) S is invertible;
- (ii) $T^*S = S^*T \iff [b(\zeta), b(\eta)] = [b^*(\zeta), b^*(\eta)] = 0;$
- (iii) $S^*S T^*T = I \iff [b(\zeta), b^*(\eta)] = \langle \zeta, \eta \rangle;$
- (iv) $ST^* = TS^*$.

A white noise operator $U \in \mathcal{L}((E), (E)^*)$ satisfies the following intertwining properties:

$$Ua(\zeta)=b(\zeta)U, \qquad Ua^*(\zeta)=b^*(\zeta)U, \qquad \zeta\in E,$$

if and only if U is of the form:

$$\begin{split} U &= C \, \exp \, \left\{ -\frac{1}{2} \Delta_{\mathrm{G}}^*(TS^{-1}) + \Lambda((S^{-1})^* - I) + \frac{1}{2} \Delta_{\mathrm{G}}(S^{-1}T) \right\} \\ &= C \, e^{-\frac{1}{2} \Delta_{\mathrm{G}}^*(TS^{-1})} \Gamma((S^{-1})^*) \, e^{\frac{1}{2} \Delta_{\mathrm{G}}(S^{-1}T)}, \end{split}$$

where $C \in \mathbb{C}$.

Proof.

 ${f 0}$ By the above two theorems, ${m U}$ is of the form

$$egin{aligned} U &= ext{wexp} \, \left\{ -rac{1}{2} \Delta_{ ext{G}}^*(TS^{-1}) + \Lambda((S^{-1})^* - I)
ight\} \diamond F \ &= ext{wexp} \, \left\{ -rac{1}{2} \Delta_{ ext{G}}^*(TS^{-1}) + \Lambda((S^{-1})^* - I) + rac{1}{2} \Delta_{ ext{G}}(S^{-1}T)
ight\} \diamond G, \end{aligned}$$

where $F,G\in\mathcal{L}((E),(E)^*)$ satisfy

$$D_\zeta^+F=0, \quad (D_\zeta^--D_{T\zeta}^+)G=0, \quad ext{for all } \zeta\in E.$$

We see from the above identity that

$$G=F\diamond ext{wexp}\,\left\{-rac{1}{2}\Delta_{ ext{G}}(S^{-1}T)
ight\}.$$

Since the right hand side contains no creation operators, we have

$$D_{\zeta}^{+}G = 0, \qquad \zeta \in E.$$
⁽¹³⁾

Then,

$$0 = (D_{\zeta}^{-} - D_{T\zeta}^{+})G = D_{\zeta}^{-}G, \qquad \zeta \in E,$$
(14)

so G is a scalar operator.

Final Remarks (1)

• We have derived a general form of U by means of a new type of a differential equation for white noise operators:

$$U = C \operatorname{wexp} \left\{ -\frac{1}{2} \Delta_{\mathrm{G}}^{*}(TS^{-1}) + \Lambda((S^{-1})^{*} - I) + \frac{1}{2} \Delta_{\mathrm{G}}(S^{-1}T) \right\}$$
$$= C e^{-\frac{1}{2} \Delta_{\mathrm{G}}^{*}(TS^{-1})} \Gamma((S^{-1})^{*}) e^{\frac{1}{2} \Delta_{\mathrm{G}}(S^{-1}T)}$$

This is the normal-ordered exponential of a quadratic function of quantum white noise (Bogoliubov Hamiltonian).

- We can derive conditions for unitarity (e.g., by using complex white noise)
- U is the composition of the generalized Fourier-Mehler and Fourier-Gauss transforms. Unitarity conditions with respect to another inner product? (Some results for *G*_{U,V}, see [Ji-Obata (2006)]).

Definition (Chung–Ji (1997))

$$\mathcal{G}_{U,V} = \Gamma(V) \, e^{\Delta_{\mathbf{G}}(U)}, \quad U \in \mathcal{L}(E,E^*), \quad V \in \mathcal{L}(E,E^*)$$

is called a generalized Fourier–Gauss transform and its adjoint operator $\mathcal{G}_{U,V}^*$ a generalized Fourier–Mehler transform.

34 / 35

We discussed

• Quantum white noise derivatives and their applications to the implementation problem for CCR.

U. C. Ji and N. Obata: A new approach to implementation problem in terms of quantum white noise derivatives, preprint, 2009.

Another applications of quantum white noise derivatives

 Hitsuda-Skorohod quantum stochastic integrals — adjoint action of derivatives U. C. Ji and N. Obata: Quantum stochastic integral representations of Fock space operators, Stochastics 81 (2009), 367–384.

Prepresentations of quantum martingales — a direct formula for the integrands U. C. Ji and N. Obata: Annihilation-derivative, creation-derivative and representation of quantum martingales, Commun. Math. Phys. 286 (2009), 751–775.